

Fig. 7. $(f_{\text{HFS}} - f_{\text{DS}})$ plotted vs atomic number at $\sin \theta / \lambda = 0.20$ and 0.70 Å^{-1} .

effects are negligible and these two models differ only in this respect. The differences are essentially zero. (The small deviation from zero of the plotted differences at low atomic number arises for the most part from small inaccuracies in the analytic fits used to compute the differences). At the higher atomic numbers the difference is always negative because the relativistic DS atomic model is more compact than the non-relativistic HFS model.

As a general rule, the inclusion of exchange has a greater effect on the scattering factors than does relativity. At the heaviest elements the two effects appear to be comparable.

Conclusions

Crystallographers now are overwhelmed with scattering factors and one might well ask which of these he should use. For the lighter elements, up to Rb⁺, the values given in *International Tables for X-ray Crystallography* (1962) are the best ones, for they have been computed from Hartree-Fock wave functions and in some cases from even more accurate wave functions. For these lighter elements, relativistic effects are trivial.

For heavier elements it is suggested that the Dirac-Slater scattering factors are to be preferred because relativistic effects are accounted for and at least an approximation for exchange has been made.

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The Crystal Structure of [18]annulene. I. X-ray Study

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The crystal structure of [18] annulene has been determined by least-squares analysis of photographic X-ray data recorded at the temperature of boiling nitrogen. The molecules occupy crystallographic centers of symmetry in a structure closely resembling that of coronene. This molecular symmetry rules out the possibility of a structure with alternate long and short C-C bonds. Final carbon coordinates have estimated standard deviations averaging less than 0.004 Å; the hydrogen atoms have not been reliably located.

Introduction

The synthesis by Sondheimer and his collaborators (Sondheimer & Wolovsky, 1959; Sondheimer & Gao-

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ni, 1960, 1961, 1962; Sondheimer, Wolovsky & Gaoni, 1960; Sondheimer, Wolovsky & Amiel, 1962; Sondheimer, 1963) of the monocyclic alternant hydrocarbons C_nH_n (n=14,16,18,20,24,30) has made available a class of compounds ('annulenes') of considerable

theoretical interest. In cooperation with Prof. Sondheimer, Dr R. Wolovsky, and Dr Y. Gaoni, we have undertaken structure analyses of two compounds in this series that are sufficiently stable to permit X-ray work. An account of the partial analysis of [14]annulene has been given by one of us (Bregman, 1962); preliminary results of the structure determination of [18]annulene were presented at the Cambridge meeting of the I.U.Cr. (Bregman & Rabinovich, 1960). The present paper reports the three-dimensional analysis of [18]annulene near 80 °K; the following paper (part II) gives the results of this study together with a discussion of the in-plane and out-of-plane molecular deformations.

Experimental

Crystals of [18]annulene, grown from slowly cooled solutions in chloroform, are thin monoclinic laths, dark brown, with occasional greenish sheen. They are elongated along [010], showing $\{001\}$ and $\{20\overline{1}\}$, with pronounced cleavage along planes parallel to [010]. Room-temperature Cu Kα photographs, which record few reflexions beyond $\sin \theta = 0.32$, yield the cell constants listed in Table 1. Systematically absent are reflexions h0l with h odd, 0k0 with k odd; the space group is $P2_1/a$. At the temperature of boiling nitrogen these crystals reflect throughout the reciprocal-lattice region accessible to Cu $K\alpha$ radiation; cell dimensions (Table 1) were determined from high-angle reflexions. The systematic absences of $P2_1/a$ were obeyed throughout the range of recorded reflexions. With two molecules of C₁₈H₁₈ per unit cell, the calculated density is 1.144 at 80 °K. The space group then demands that the molecule be centrosymmetric, unless the structure is disordered. This last contingency has been ruled out by refinement of an ordered model to a structure yielding acceptable agreement between F_o and F_c and showing no anomalies in the thermal parameters (see part II). [18] Annulene cannot, therefore, be formulated as a molecule with alternate single and double bonds.

Examination of a large number of crystals revealed the frequent occurrence of twinning across (001), made possible by the near equality (within 1.2%) of c^* and $4a^*\cos\beta^*$. The twinned crystals give, at room temperature, X-ray photographs that appear to have been produced by crystals having a doubled c axis and showing non-space-group extinctions for h even, L (=2l) odd. At low temperature high-angle spots are found to be split, suggesting that the phenomenon observed is one of twinning rather than of polymorphism or

super-lattice formation. This conclusion is confirmed by the fact that the positions and intensities of all spots appearing on photographs of a twinned crystal can be exactly simulated by the superposition in mirror orientation of pairs of reciprocal-lattice nets derived from photographs of untwinned crystals. This superposition brings into near coincidence spots due to reflexions h, k, l and h, k, $-\frac{1}{2}(h+2l)$.

Reflexions from an untwinned crystal were recorded with Cu $K\alpha$ radiation by the equi-inclination Weissenberg technique in levels perpendicular to [010] up to k=3. The crystal was held near the temperature of boiling nitrogen by the device described by Hirshfeld & Schmidt (1956). From a crystal trimmed with a razor blade 0kl intensity data were collected that enabled us to place I(hkl) (k=0,1,2,3) on a common scale. However, the quality of cut crystals was inadequate to permit intensity photographs about another axis. The numbers of reflexions, observed plus unobserved, were 183, 356, 335, and 284, respectively, for k=0,1,2,3 and 3. Ten 04l reflexions were also included in the refinement process.

The intensities were estimated visually and corrected in the usual way; corrections for spot-shape extension, due to Phillips (1956), were applied to the data from upper-level photographs.

Solution and refinement

Comparison of the cell dimensions of [18]annulene and of coronene (Table 1) suggests a similarity in molecular shape and packing arrangement; since the strong high-order h0l reflexions (207; $40\overline{5}$; 602; $10,0,\overline{3}$; 14,0,1; $14,0,\overline{8}$; $16,0,\overline{1}$), used by Robertson & White (1945) for the solution of the latter structure, were, with the exception of 14,0,1, outstandingly strong in [18]annulene too, the x, z parameters of the peripheral carbon atoms of coronene were inserted into F(h0l) structure-factor and least-squares calculations. The correctness of this model quickly became evident.

Refinement was accomplished on WEIZAC by a least-squares program modelled after the procedure described by Rossmann, Jacobson, Hirshfeld & Lipscomb (1959). This program seeks to minimize the discrepancy function

$$r = \sum w(K^2 F_0^2 - F_0^2)^2 / \sum wK^4 F_0^4$$
 (1)

with respect to the scale factor K^2 and the positional and thermal parameters of the atoms in the asymmetric unit. It uses the linear diagonal approximation for

Table 1. Crystal data for [18]annulene and coronene

	a (Å)	b (Å)	c (Å)	β	Z	Space group
$C_{18}H_{18}$ at $80^{\circ}K$	14·889 + 0·004	4·800 + 0·002	10·235 + 0·003	111·60° + 0·14°	2	$P2_1/a$
C ₁₈ H ₁₈ , room temperature	15.33	4.88	10.27	111·8°	2	$P2_1/a$
Coronene, room temperature	16.10	4.695	10.15	110·8°	2	$P2_1/a$

the computation of the shifts in all parameters, except that the scale factor K^2 and an average, isotropic, thermal parameter are adjusted separately by solution of an appropriate 2×2 matrix (Cruickshank, 1961).

For each reflexion, the weighting factor w_{hkl} was obtained by the summation of weights assigned to the individual spots on the several films on which the reflexion was recorded. These individual weights were a function of the film blackening and of the factors required for translation of the spot readings into scaled values of F_o^2 . This procedure yielded weights that varied roughly as F_o^{-4} for most reflexions, with decreased weights for the weak and the very strong reflexions. The summations in equation (1) comprise all except those unobserved reflexions for which $|F_c| < KF_t$, where F_t is a threshold value estimated for each unobserved reflexion. Scattering-factor curves were taken from Berghuis, Haanappel, Potters, Loopstra, Mac-Gillavry & Veenendaal (1955) for carbon and from McWeeny (1951) for hydrogen.

Refinement of the (010) projection proceeded satisfactorily, initially with isotropic carbon atoms only. Next, hydrogen coordinates derived from a difference synthesis $\Delta \varrho(x, z)$ were inserted into the least-squares

calculations, and the refinement of the carbon parameters was continued for several additional cycles.

In the next stage y parameters of all atoms were derived for a hypothetical planar molecule; this model was tested by structure-factor calculations for F(h11) and found to give reasonable agreement.

Because of the diagonal approximation used in the least-squares program, it was desirable to refer the structure to nearly orthogonal axes. The coincidence in cell dimensions responsible for the twinning (see above) provided such a coordinate system, based on a unit cell four times the size of the primitive cell. The axes of this cell, having a monoclinic angle of 90·26°, are obtained from the conventional axes by the transformation whose matrix is

$$\left(\begin{array}{ccc} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 1 & 0 & 4 \end{array}\right)$$

In the first four three-dimensional refinement cycles, only the scale factor K^2 , the y parameters, and the isotropic thermal parameters (B) of the carbon atoms were varied; next all the carbon parameters were re-

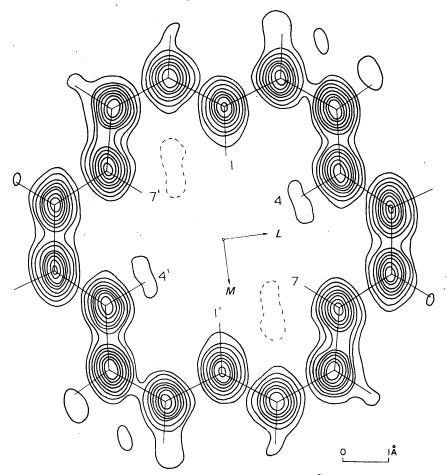


Fig. 1. Electron density in mean molecular plane. Contours at intervals of 1 e.Å-3, zero contour broken. Molecular outline shows projected atomic positions based on least-squares refinement.

Table 2. Atomic coordinates

	Refer	red to crystal ax	es	Reie	terred to molecular axes			
Atom	x	y	Z	L(Å)	M(Å)	N(Å)		
C(1)	-0.1147	-0.4574	0.0430	0.424	-2.881	0.061		
C(2)	-0.1068	-0.5089	0.1807	1.742	-3.323	-0.015		
C(3)	-0.0441	-0.3627	0.2988	2.864	− 2·467	-0.058		
C(4)	0.0151	-0.1483	0.2903	2.775	−1 ·091	-0.053		
C(5)	0.0836	-0.0082	0.3995	3.807	-0.178	0.010		
C(6)	0.1430	0.2077	0.3781	3.594	1.236	0.041		
C(7)	0.1416	0.2877	0.2474	2.343	1.828	0.085		
C(8)	0.1885	0.5084	0.2154	2.034	3.162	-0.013		
C(9)	0.1771	0.5884	0.0770	0.715	3.675	-0.064		
H(1)	-0.072	-0.278	0.031	0.31	−1·79	0.00		
H(2)	-0.143	-0.668	0.197	1.90	-4.30	0.04		
H(3)	-0.043	-0.427	0.390	3.74	- 2 ⋅91	-0.06		
H(4)	0.006	 0·078	0.195	1.86	- 0·69	-0.14		
H(5)	0.095	- 0·054	0.499	4.76	 0·47	0.05		
H(6)	0.187	0.320	0.464	4.41	1.87	-0.02		
H(7)	0.107	0.168	0.174	1.64	1.24	0.25		
H(8)	0.227	0.626	0.297	2.81	3⋅77	-0.15		
H(9)	0.211	0.749	0.060	0.55	4.63	-0.16		

fined. After several further cycles anisotropic carbon thermal parameters β_{ij} (i, j=1, 2, 3) were introduced (Cruickshank, 1956), and only when the nine parameters per carbon atom had reached essentially stationary values were the positional and isotropic thermal parameters of the hydrogen atoms allowed to vary. After 28 cycles the discrepancy factor r had reached 0.0372.

At this stage bond lengths and angles were calculated, as well as the equation of the best plane through the carbon skeleton and the deviations of the carbon and hydrogen atoms from this plane. The carbon skeleton conformed closely to symmetry 3; on the other hand, the displacements of hydrogen atoms from the mean plane did not follow a similar pattern. In particular, the position of H(1) appeared, on the criterion of molecular symmetry 3, to be in serious error. A three-dimensional difference synthesis based on $F_o - F'_c$ (with the sign of F_c , including hydrogen contributions, attached to F_o ; carbon contributions alone for F'_c) was computed in the neighborhood of H(1) and H(4). While for the latter the electron-density maximum coincided within 0.02 Å with the position derived by the leastsquares procedure, the electron density near H(1) indicated a position 0.11 Å above, rather than 0.14 Å below, the mean plane of the carbon atoms. This more plausible location for H(1) was inserted into further least-squares calculations, but these returned the atom to its former position.

The entire set of intensity data was re-checked for human errors, and minor corrections were made in the intensities of about 100 reflexions, especially at high angles. About 60 further reflexions were assigned zero weights because of high background scattering or because of suspected interference by the beam catcher. New hydrogen coordinates were postulated, based on the last set of carbon coordinates and the assumption that each C-H bond, of length 1.00 Å, lay in the plane defined by the three adjacent carbon atoms and bisected the C-C-C angle. The positional and thermal parameters of the carbon atoms were refined for another five cycles, after which the hydrogen parameters were again allowed to change.

The final hydrogen positions did not show any significant improvement, though the anomaly in H(1) was smaller than before. A new three-dimensional difference synthesis, calculated in the same way as the previous one, showed maxima at positions differing by up to 0.14 Å from the hydrogen positions given by least-squares. All hydrogen atoms showed marked anisotropy; the electron-density distributions resembled

Table 3. Direction cosines of the molecular inertial axes with respect to the crystal axes

	a	b	\boldsymbol{c}
L	-0.0393	0.0266	0.9433
M	0.6593	0.7517	-0.2370
N	0.7509	-0.6587	-0.2325

Table 4. Estimated standard deviations (average) of atomic coordinates and thermal tensor components, referred to the edges of a nearly orthogonal cell having a = 14.889, b = 4.800, c' = 38.059 Å, $\beta' = 90.26^{\circ}$ (This is the cell used for the diagonal least-squares refinement)

Carbon: $a\sigma(x') = 0.0024 \text{ Å}$ $\sigma(U'_{11}) = 0.0010 \text{ Å}^2$ $\sigma(U'_{12}) = 0.0013 \text{ Å}^2$	$b\sigma(y) = 0.004 \text{ Å}$ $\sigma(U'_{22}) = 0.0026 \text{ Å}^2$ $\sigma(U'_{23}) = 0.0014 \text{ Å}^2$	$c'\sigma(z') = 0.0023 \text{ Å}$ $\sigma(U'_{33}) = 0.0010 \text{ Å}^2$ $\sigma(U'_{13}) = 0.0008 \text{ Å}^2$
Hydrogen: $a\sigma(x') = 0.030 \text{ Å}$	$b\sigma(y) = 0.038 \text{ Å}$ $\sigma(u^2) = 0.008 \text{ Å}^2$	$c'\sigma(z') = 0.027 \text{ Å}$

Table 5. Observed and calculated structure factors

Those marked with an asterisk were eliminated from the final refinement because of high background scattering or because of possible interference by the backstop.

<u>h</u> <u>k</u> <u>l</u>	100 <u>F</u> _o 100 <u>F</u> _c	<u>h</u> <u>k</u> <u>1</u> 100 <u>F</u> ₀ 100 <u>F</u> _c	<u>h</u> <u>k</u> <u>1</u> 100 <u>F</u> _o	100 <u>F</u> c	<u>h</u> <u>k</u> <u>1</u>	100 F ₀ 100 F _c	<u>h</u> <u>k</u> <u>1</u>	100 <u>F</u> 10
18 0 -1* -2 -3 -4 -5 -6 -7 -8*	1325 1589 713 713 540 -480 458 4,6 342 -376 223 234 531 492	-11° (41 -785 6 0 10° 102 -324 9 880 771 8 390 334 7 1812 -1838 6 15:11 -1449 5 16:10° -1506	-13 248 -12 663 -11 599 -10 219 -9 505 -8 <231 -7 810	-259 -630 -592 -220 481 163 71 ¹	-3; -10 -9 -6 -7 -6 -5	2-30 2677 1525 -1606 700 715 671 617 380 -410 1372 1319 278 219	-4 -3 -2 -1 0 1	395 3 1000 9 1137 12 1532 -15 2831 30 964 -9 387 4 1248 13
16 0 2* 1 0 -1 -2	946 -1162 480 684 967 -971 313 268 2386 2307 763 -743	4 10°6 1097 5 <188 -1192 8 4057 -4,381 1 486 304 0 1790 1792	-6 <231 -5 569 -1 360 -3 76; -2 600 -1 271 0 591 1 481	-90 818 265 684 602 229 561	-3 -2 -2 0 1	832 -836 481 -519 197 81 12-7 1367 14-17 1512 <2.2 142 <212 209	3 5 6 7 8	286 271 2 533 5 423 -3 <219 1 350 -5
-3 -4 -5 -6 -7 -8 -9	513 469 <203 214 <249 69 603 -528 1025 1103 342 396 547 -582 844 -871	-1 975 854 -7 2740 1540 -3 1540 1559 -1 994 859 -5 551 -562 -6 <48t -7 954 938 -8 <219 93 -8 <219 -9 <226 -113	2	-401 582 -668 -762 72 -53 607	3 5 6 7 8	1160 -1202 702 -699 948 1028 <212 23 248 266 635 655 <205 -305 <170 -189	10 11 12 13 -18 1 3 -17 -16	651 -5. <2231. 467 4. 4014. <197 1. <1881.
14 0 4 3 2 1 0	<214 33 417 343 1785 -1723 611 -650 618 -682	-10 <226 -8 -11 626 587 -12 <175 -59 -13 379 313 4 0 10 553 526 9 472 472	-12 497 -11 <242 -10 724 -9 682 -8 <212 -7 334	-375 -78 -630 721 144 351 362 -859	10 11 -18 1 6 -17 -16 -15 -14	<107 146 595 502 <197 -81 395 -306 174 340	-15 -14 -23 -12 -11 -10	553 -4 306 -3 505 -4
-2 -3 -4 -5 -6 -7 -8	975 -982 962 972 844 849 707 695 1098 -1182 611 630 3755 3377	7 1696 1756 6 1310 1266 5 1440 -1481 4 1696 1613	-6 372 -5 854 -4 262 -3 289 -2 510 -1 <204 0 <197	-859 -266 -280 -574 239 77 780 162	-13 -12 -11 -10 -9 -8	344 -260 927 964 <219 131 <212 255 964 -987	-9 -8 -7 -6 -5 -4 -3 -2	2197 1. 365 -3 358 2 1672 -15; 620 -6 255 -2 1379 -12; 212 2 271 -2
-9 -10 -11 -12 12 0 6 5	3755 3377 349 -441 822 -762 618 620 <79 -101 240 273 2214 -70 1034 -925	0 5328 -5116 -1 7125 -6315 -2 4533 -4259 -3 2038 2047 -4 540 -534	0 <197 1 795 2 <223 3 205 4 818 5 547 -17 1 9 818 -16 <146	236 -733 -508 893 208	-7 -6 -5 -4 -3 -2 -1	2022 -2059 <181 133 1073 1058 <176 -109 780 705 1313 -1368 181 -325 1445 -1501	-1 0 1 2 3	1904 199 <125 119663 -796 2481 -2596 1795 -186 1532 -1696
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-9 -10 -11 -12 10 0 7 6 5	815 826 707 766 298 295 <186 153 1034 1123 801 911	3 1206 1201 2 1529 1464 1 5952 -5890 0 2373 2123 -1 1382 1229 -2 5100 -5498	-1 <219 0 1430 1 334 2 <212 3 849 4 <231	-577 -75 -1522 -340 1 6 4 6 6 9 145	-18 1 5 -17 -16 -15 -14 -13 -12	730 727 408 -359 760 760 <204 -27 <219 207 489 545 380 356 <212 -10	-15 -14 -13 -12 -11 -10	561 -5 1058 10
3 2 1 0 -1 -2	<226 35 263 235 1170 1251 1550 1548 641 667 547 563 641 -621	-3 3369 3395 -4 916 -935 -5 291 -3 6 -6 1951 1855 -7 1288 -1230 -8 1513 -1520 -9 858 -849	5 <214 6 1921 7 <153 8 <51 -18 1 8 <94 -17 76 -16 898	156 -942 31 -68 211 736 8 6 0	-11 -10 -9 -8 -7 -6	<212 -10 787 -141 752 -759 226 -261 1642 -1754 714 -696 286 -244	-9 -8 -7 -6 -5 -4 -3	708 -6 4241 -39 1563 12 1262 -11 1664 -15
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7 5 1, 3 2	<234 -198 <226 80 707 646 1513 -1414 1208 -1042	2 4314 -4635 1* 916 945 -6 1 13* 262 -353 -13 1 12* 278 -384 -12 139 244 -11 255 259	-2 1795 -1 2342 0 635 1 511 2 313 3 1496	-521 379 1875 2588 653 -527 -298 -1616	9 10 11 12 13 -18 1 4	<186 46 561 641 132 122	11 12 13 14 15 16 *	760 -7 591 -5 <204 <223 - 481 3 <109 -1 416 5 927 -9
-1 -2 -3 -4 -5	480 476 1607 1511 946 826 291 139 931 875 1201 1119	-10 240 317 -9 481 9 -8 321 270 -7 474 -454 -6 488 70 -5 787 -755 -4 181 -156	5 54 0 6 730 7 423 8 328 9 459	1178 -548 747 -418 -395 469 -576 -205	-17 -16 -15 -14 -13 -12 -11	606 598 387 -215 627 -674 752 -809 584 -556 431 -390 803 760	-17 -16 -15 -14 -13 -12	<269 1 627 5 905 -8 599 5
-6 -7 -8 -9 -10 -11	1819 1821 924 896 226 215 234 -238 1477 -1817 <249 -52 249 354	-3 452 -477 -2 372 -337 -1 299 219 0 <117 -62 -15 1 11 511 665 -14 <125 38	-16 1 7 540 -17 181 -16 986 -15 1174 -14 <212 -13 <219 -12 406	-205 -976 -1022 82 -135 -414	-10 -9 -8 -7 -6	803 760 395 -325 1110 1066 <176 -110 1825 1790 810 783 1460 1539	-11 -10 -9 -8 -7 -6	1183 10 293 2 2183 -20 1956 -16 547 -4

Table 5 (cont.)

<u>h</u> <u>k</u> <u>l</u>	100 <u>F</u> o	100 <u>F</u> c		h s	1	100 F _o	100 <u>F</u> c	h k j	<u>.</u>	100 <u>F</u> o	100 F _c	<u>h</u> k	ī	100 <u>F</u> o	100 F _e	<u>h</u> k		100 <u>F</u> 0	100 <u>F</u> c
-5 -4	4 0 43 6269	3792 5962 8 0 41		-2 -1 0 1		14 0 7 <197	-1592 22	-12 -11 -10		616 <2 0 5 7 0 1	556 76 658	-7 -6		18 04 89 0 <135 3154	1766 8 0 5 •54	-14 3 -12 -11	9	677 265 434 236	755 -209 -429 -174
-3 -2 * -1 *	7918 1313 4561	8041 1233 -4088		0		<191 357 582	-214 269 474	-4		1211	658 1087 -498	-5 -4 -3		<135 3154 941	-54 -3343 912	-10		324	
0 .	350 4670	429 - 38 0 0		3		607 155	-511 119	-8 -7 -6		1033 326 855	1009 -290 903	-3 -2 -1		340 354	281 344	-9 -8		1325 972 1606	1270 959 1637
2 3 4	3773 1539	4027 -1605 2246		5	4	326 629	-346 552 -354 -18	-5 -4 -3		389 1160	-340 -340	5 0		<7145 383 858	894 -355	-7 -6 -5 -4		1605 < ?79 <228	1270 959 1637 -199 -151 -819
5 6	2240 176 <153 1547	-141 -66		-17 2 -16	8*	<57 <92 411	388	-3 -2		6 0 ₿ <162	559 137 1540	3		1062	1350 -36 1026	-4 -3 -2		868 <279	
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Table 5 (cont.)

5 <u>k</u> <u>1</u>	100 <u>F</u> o	100 £;	<u> </u>	ķ	ī	100 <u>P</u> o	100 <u>F</u> _c
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20 112 ° 16 3 3 1-15 14 14 14 14 14 14 14 14 14 14 14 14 14	150 5 / 2 6 / 3 2 / 5 3 / 5 4 / 7 7 / 8 3 / 5 4 / 6 7 / 6 6 / 6 6 / 6 6 / 6 6 / 6	8 179 -710 -146 -247 -246 -511 -46 -517 -492 177 3-492 177 3-492 177 3-492 177 3-562	4 5 6 6 7 8 9 9 10 11 12 12 13 12 12 12 12 12 15 6 5 5 4 6 6 7 6 7 6 6 5 6 6 7 6 7 6 6 7 6 7 6 7	3	0	510 1347 1355 2180 4243 425 426 426 426 426 426 426 426 426 426 426	-779 1293 1296 2247 2247 231 -1048 -149 151 151 147 80 678 -189 -195 -195 -254 -254 -254 -254
5 6 7 8 9 10 21 12 13 13 2-16 3	500 214	-562 -776 - 022 - 032 - 103 - 103 - 133 - 19 - 216 - 235	3 2 0	l.	* 98 7 6 5 4 3 2 1 0	2377 840 1221 <145 348 551 268 1379 <283 642 <276 <276	759 1167 -14 470 -679 -183 1735 -178 -667 8

prolate spheroids with their long axes nearly normal to the molecular plane. Final values of r and $R = \Sigma |KF_o - |F_c||/\Sigma KF_o$ were 0.0289 and 0.076, respectively; unobserved reflexions for which $KF_t < |F_c|$ were included in the sums.

Observed and calculated structure factors are compared in Table 5.

An electron-density section in the mean molecular plane is shown in Fig. 1.

The final coordinates of all atoms in the asymmetric unit, derived from the last cycle of least-squares refinement, are listed in Table 2 with respect both to the crystal axes and to a set of orthogonal molecular axes. The latter (Table 3) were chosen to coincide with the principal molecular moments of inertia, calculated for the carbon atoms only, by the same computer program that yielded the equation of the mean plane according

to the method of Schomaker, Waser, Marsh & Bergman (1959). Because the molecule has very nearly $\bar{3}$ symmetry (see part II) the moment of inertia is almost axially symmetric; thus the directions of the in-plane inertial axes L and M have slight physical significance.

Table 4 gives the average standard deviations computed by the diagonal approximation:

$$\sigma^{2}(u_{i}) = \sum w(K^{2}F_{o}^{2} - F_{c}^{2})^{2}/(n-s)A_{ii}$$
 (2)

where n=911 is the number of reflexions, 883 observed plus 28 threshold, included in the summations of the last refinement cycle, s=118 is the number of parameters u_i adjusted, and $A_{ii} = \sum w(\partial F_c^2/\partial u_i)^2$ is the diagonal element in the normal equation corresponding to the parameter u_i . The approximate validity of the diagonal approximation follows from the near orthogonality of the axes and the large number of reflexions observed.

Fig. 2 shows the final values of bond lengths and angles and the distances of all atoms from the mean plane of the carbon atoms. This plane has the equation

$$11 \cdot 1795x - 3 \cdot 1616y - 2 \cdot 3798z = 0$$

in which the left side gives the distance from the plane, in \mathring{A} , of a point having fractional coordinates (x, y, z). The results of this analysis are discussed in part II.

We are grateful to Professor Franz Sondheimer and Dr Reuven Wolovsky, who brought this problem to our attention and supplied the crystals for this investigation. It is a privilege to acknowledge our indebtedness to the late Professor I. Fankuchen, whose suggestion led to the proper diagnosis of the twinning phenomenon.

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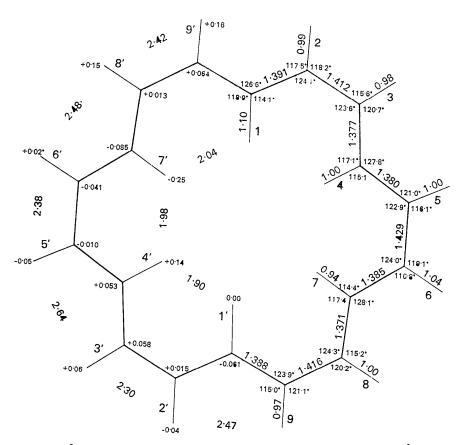


Fig. 2. Interatomic distances (Å) and bond angles. Figures near atomic positions give displacements (Å) from mean molecular plane.

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